

MASS DISTRIBUTION OF SHOWER METEORS

by

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Progress Report No. 1 (May 1 - October 31, 1966)

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THE /UNIVERSITY OF ADELAIDE

DEPARTMENT OF PHYSICS 3

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1. <u>INTRODUCTION</u>

It has been shown (Elford, 1965) that the diurnal variation in hourly rate of radar echoes from a meteor shower depends on the distribution of the number of meteors as a function of mass. If the parameters of the radar system are known with sufficient precision the theoretical echo rate for a point source radiant of unit strength can be determined for all positions of the radiant in elevation and azimith. The result of such a calculation is a description of the response of the system to a radiant in any position in the sky and is appropriately called the 'response function' of the system. The form of the response function will vary according to the particular choice of the mass distribution. Hence a comparison of the observed and predicted radar echo rates for a particular meteor shower can lead to a determination of the mass distribution.

The procedure outlined above is being applied to the results of precise meteor echo rates measured at Christchurch (43.5°S) and Ottawa (45°N). The success of the analysis depends on accurately matching observed and predicted echo-rate versus time curves. In the first part of this project emphasis has been given to determining theoretical echo rate v. time plots for various mass distributions and for a range of values of declination of a point radiant. For this purpose the response functions of the radar systems at Christchurch and Ottawa have been recalculated in order to take into account overdense echoes

and the range interval used in recording the radar data. These calculations are described in Section 2.

In order to reduce the number of graphs of the diurnal variation in echo rate a method of plotting contours of equal echo rate on a declination-time plot has been developed. The results are described in Section 3. The effect of a change in the mass distribution law on the rate plots is also discussed in Section 3.

The published data on radar echo rates determined at Christchurch give the total echo rate (shower plus sporadic) at any hour, and in order to determine the shower echo rate alone it is necessary to remove the sporadic background. This procedure is discussed in Section 4.

2. RESPONSE FUNCTIONS

The response functions for the Ottawa and Christchurch systems have been calculated according to the method developed by Elford (1964), but modified to include overdense trails, any form of the distribution of ionization along a trail and any range interval. The calculations have been carried out with the aid of a CDC 6400 computer and the program is given in Appendix I. For the purpose of machine calculation it was found necessary to smooth the polar diagram of the Christchurch antenna. The smoothed polar diagram data are given in Appendix II. The polar diagram for the Ottawa antenna was taken as that for crossed half-wave dipoles 0.4 wavelengths above a reflecting screen. The theoretical expression is given in Appendix II.

The incident flux of meteors whose radiants lie within an element of the celestial sphere of solid angle $d\Omega$ and which produce zenithal line densities greater than q_z can be written as $N(q_z)d\Omega$. This is the flux across a plane normal to the meteor paths. It is assumed that the cumulative distribution function $N(q_z)$ can be represented by a simple

power law of the form $N(q_z) = Kq_z^c$. Typical response functions for the Ottawa and Christchurch systems are given in Figures 1 and 2 for several values of the exponent c. It can be seen that the radiant elevation at which these systems have their peak response depends on the value of c. In the case of the Christchurch system the minor low angle lobe in the antenna polar diagram produces a very marked variation of the position of the peak response as the exponent c is changed.

The Ottawa system records meteor echoes out to a maximum range of 360 km and it is of interest to examine the effect of this range restriction on the response function. The result is shown in Figure 3. It is clear that a reduction of the maximum range from 1,000 to 360 km has only a marginal effect on the response function for the Ottawa system. The effect of a similar range restriction on the Christchurch system has also been examined and the result is shown in Figure 4. In this case the restriction in range would have had a significant effect on the response function for radiant elevations in excess of 55°. The actual rate data obtained at Christchurch had no range restriction.

3. DIURNAL VARIATION IN ECHO RATE

As a radiant moves across the sky the elevation, $oldsymbol{ heta}$, from a given station can be determined from the following expression

 $\sin\theta = \cos\frac{\pi}{12}t \cos\delta \cos L + \sin\delta \sin L$, where L is the latitude of the station, δ is the declination of the radiant and t is the time in hours after transit. Hence for a given station and a given radiant the variation in the theoretical echo rate as a function of time can be determined by calculating the radiant elevation at any given time and noting the value of the response function at this elevation.

The theoretical echo rate of meteors from point radiants as a function of time has been computed for the Ottawa and Christchurch systems. The calculations were carried out at intervals of 5 degrees in radiant declination and for five values of the exponent c in the mass distribution law, viz. - 0.6, - 0.8, - 1.0, - 1.2, - 1.4. For convenience the results are presented as contours of equal rate on a declination-time plot. The results for the Ottawa system are presented in Figures 5 - 9 and for the Christchurch system in Figures 10 - 14. The contours range from zero to a maximum rate of ten units. The approximate positions of the contours of zero rate are indicated by dashes and the positions of the contours of maximum rate by plus symbols.

The rate contour diagrams for the Ottawa system show that radiants with declinations higher than $+ 10^{\circ}$ give a rate-time variation with two maxima spaced equally either side of the time of transit. For radiants with declinations between $+ 35^{\circ}$ and $+ 55^{\circ}$ the echo rate goes to zero for about an hour before and an hour after transit. The main effect of a change in the value of c from - 0.6 to - 1.4 is to expand the region occupied by the echo rate 'well'. Hence at declinations higher than $+ 10^{\circ}$ the time of occurrence of the peak rate can change by as much as one hour when c is changed from - 0.6 to - 1.4.

The general features of the rate contour diagrams for the system at Christchurch are similar to those for the Ottawa system but the effect of a change in the value of c is more marked. For c = -0.6 (Figure 10) a reduction in echo rate near transit only occurs for radiants lying between declinations -25° and -60° and the maximum duration of the decrease in echo rate is three hours. In Figure 14 the contours of constant echo rate are drawn for c = -1.4 and from

this diagram it can be seen that all radiants south of declination + 5° give rate time diagrams with a minimum at the time of transit and that for radiants lying between declinations 0 and - 75° the duration of the decrease in echo rate ranges from 4 to 14 hours.

Typical cross-sections of an echo rate contour diagram are given in Figure 15 where rate-time curves are plotted for the Christchurch system at intervals of 5° in radiant declination and for c = -1.0. These curves show that for a radiant with a medium southern declination the observed echo rate as a function of time can exhibit quite complicated behaviour.

We have already seen that the whole iso-rate diagram is altered by a change in the mass distribution law. A close examination shows that the effect of these changes is most marked for radiants whose zenith distance at transit is less than 30°. The extent of the dependence of the echo rate on the mass distribution law is illustrated in Figures 16 and 17 where normalized diurnal rate curves for the Ottawa and Christchurch systems are plotted for three values of c and for two values of declination. In the case of the Ottawa system the diurnal rate curves tend to become narrower as the value of c becomes more negative. For the Christchurch system changes in the values of c cause only marginal changes in the form of the rate curves for low declination radiants. However for radiants at high declination a change of c from - 0.6 to - 1.4 completely alters the form of the diurnal rate curve and even a change of 20% in c produces significant changes in the rate-time plot.

4. OBSERVED DIURNAL RATE VARIATIONS

Detailed accounts have been published of radar echo counts of meteors carried out at Ottawa and Christchurch. In order to determine the diurnal variation in shower activity an estimate of the sporadic

activity must be subtracted from the total echo rate. This is being carried out in the following way. The published data have been replotted in terms of the day-to-day variation in total echo rate during one year for each hour of local time. The times of occurrence of known showers are marked on the annual plots and a base rate is drawn by eye. Any inconsistencies in successive hourly plots are noted and appropriate corrections made. The base rate is subtracted from the total echo rate and the diurnal shower activity plotted for several days near the peak of each shower. The chief uncertainties occur when showers overlap in time and in general it is not possible to deduce diurnal rate curves in these cases.

Results of radar echo counts of Geminid meteors observed with the Ottawa system have been published by McIntosh (1966). The rates of the short duration echoes as a function of radiant elevation angle were compared with that predicted using c = -0.5 and - 2.0. Neither of these rate curves fit the observations. The width of the observed curves is indicative of a value of c close to - 2.0 but the position of the peak suggests c = -0.5. Thus a simple inverse power mass distribution is not adequate to describe the observed diurnal rate variations of Geminid echoes. A better approximation would be obtained by replacing the single power term by the first three terms of a power series whose coefficients are to be determined by the comparison of predicted and observed rate curves. A program to give the appropriate response functions is being written.

APPENDIX I

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APPENDIX II Antenna polar diagram data

(a) Christchurch: Crossed folded dipoles above a ground screen.

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(b) Ottawa: Crossed dipoles 0.4λ above a ground screen.

Theoretical polar diagram

$$G(\theta) = \left[\frac{\cos^2(\frac{\pi}{2}\cos\theta)}{\sin^2\theta} + 1 \right] \sin^2(0.8\pi \sin\theta)$$

REFERENCES

- Elford W. G. Calculation of the response function of the Harvard Radio Meteor Project radar system. Harvard-Smithsonian Radio Meteor Project Res. Report. No. 8. (1964).
- Elford W. G. The incidence of meteors on the earth derived from radio observations. Accepted for publication in the Smithsonian Contributions to Astrophysics.
- McIntosh B. A. The determination of meteor mass distribution from radar echo counts. Can. J. Physics. 44, 2729 (1966).

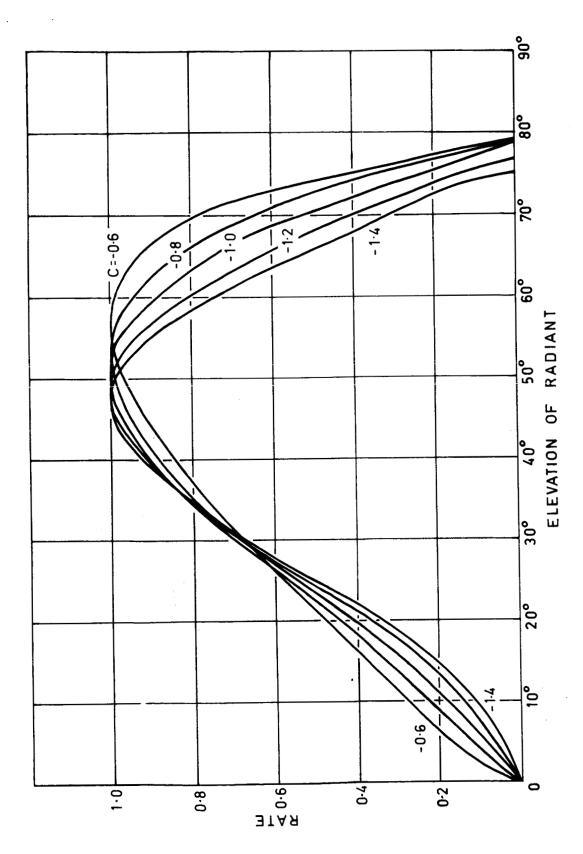


Figure 1. Response functions for the 33 mc/s system at Ottawa, calculated for several values of the exponent C in the distribution law. Meteors detected between the range limits 100-360 km.

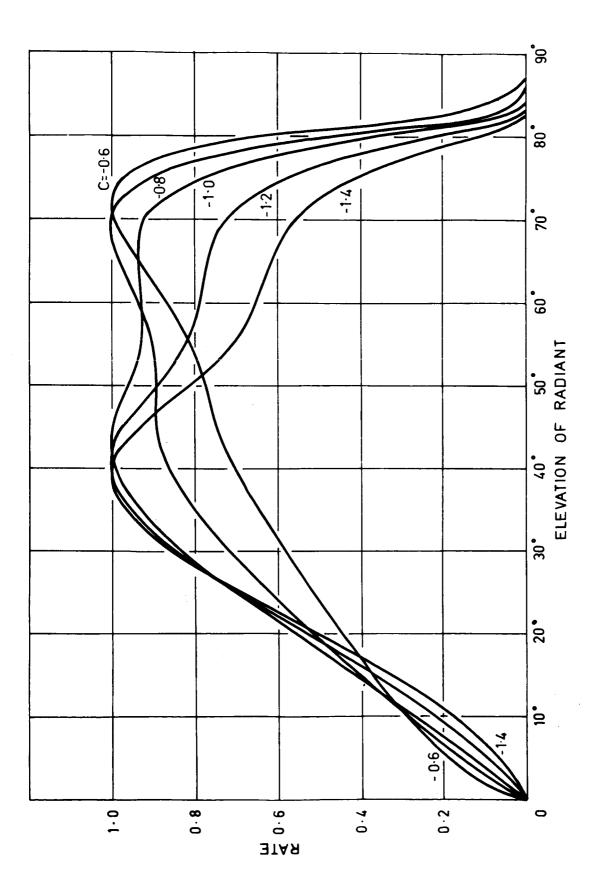


Figure 2. Response functions for the 70 mc/s system at Christchurch, calculated for several values of the exponent C in the distribution law. Meteors detected to a maximum range of 1,000 km.

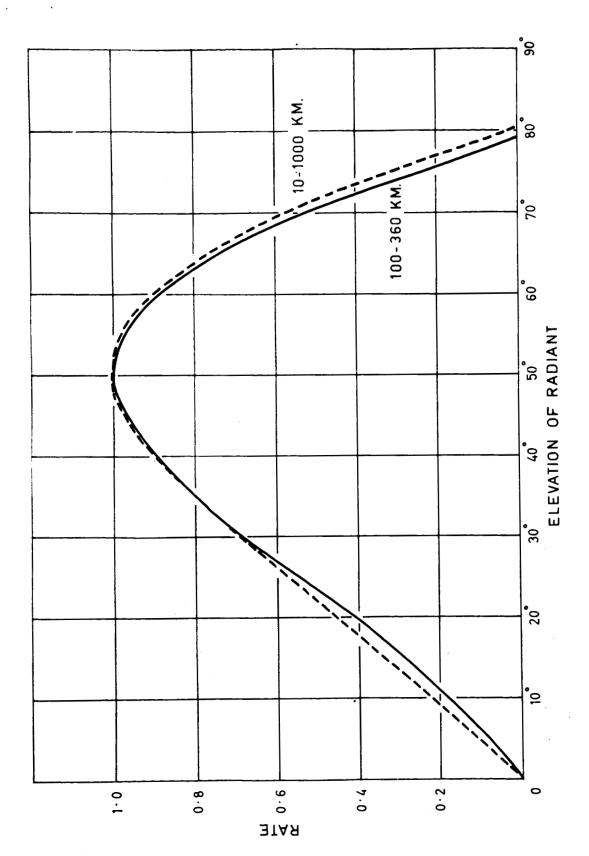


Figure 3. Normalized Ottawa response functions for two range limits (C = -1.0).

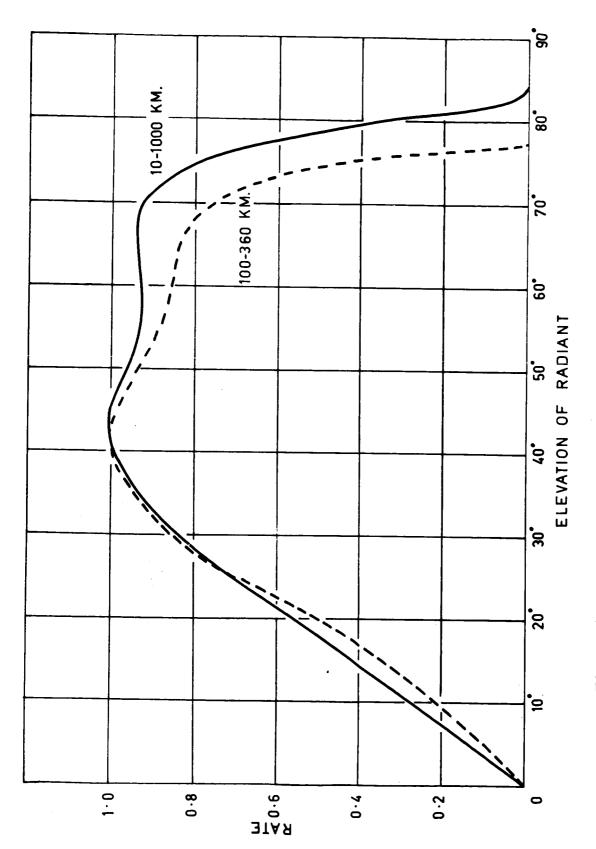
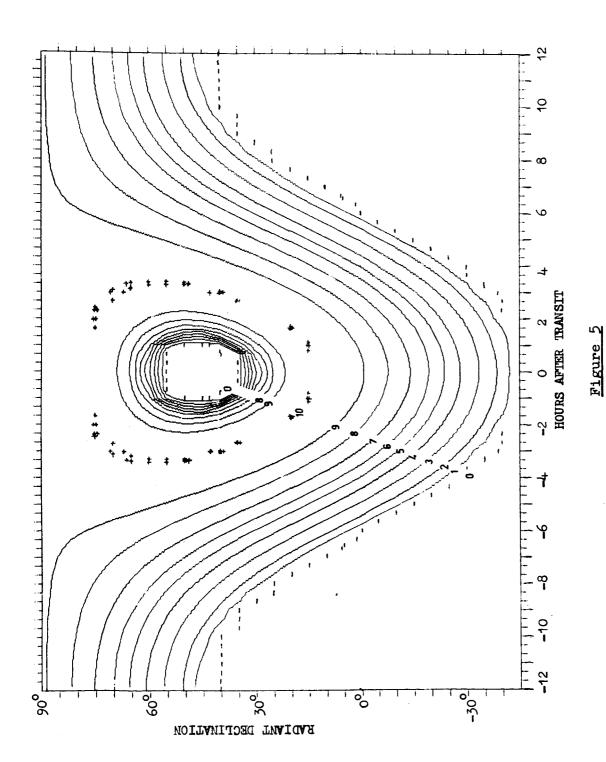
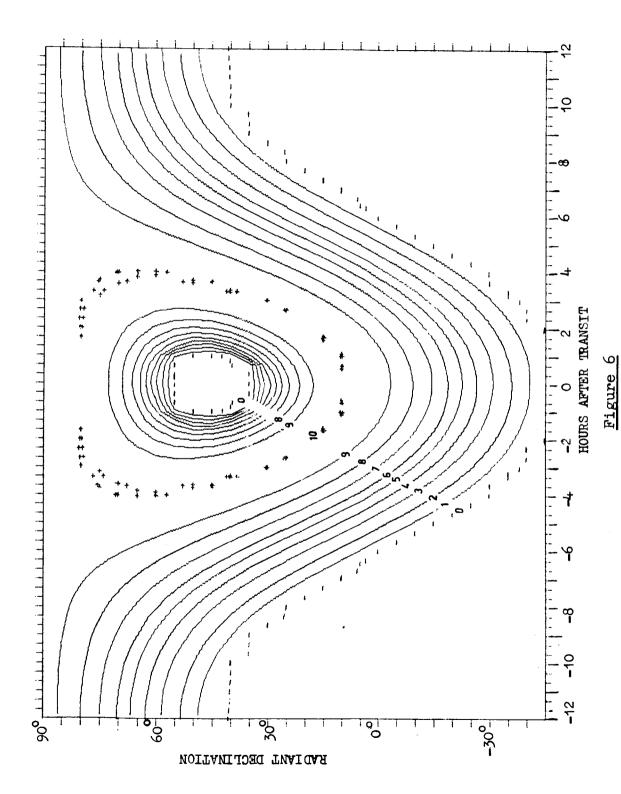


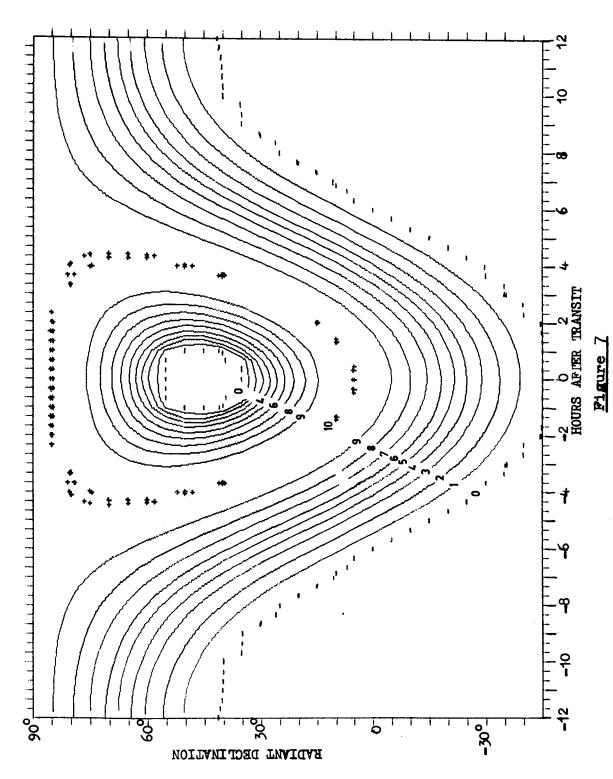
Figure 4. Normalized Christchurch response functions for two range limits (C = -1.0).



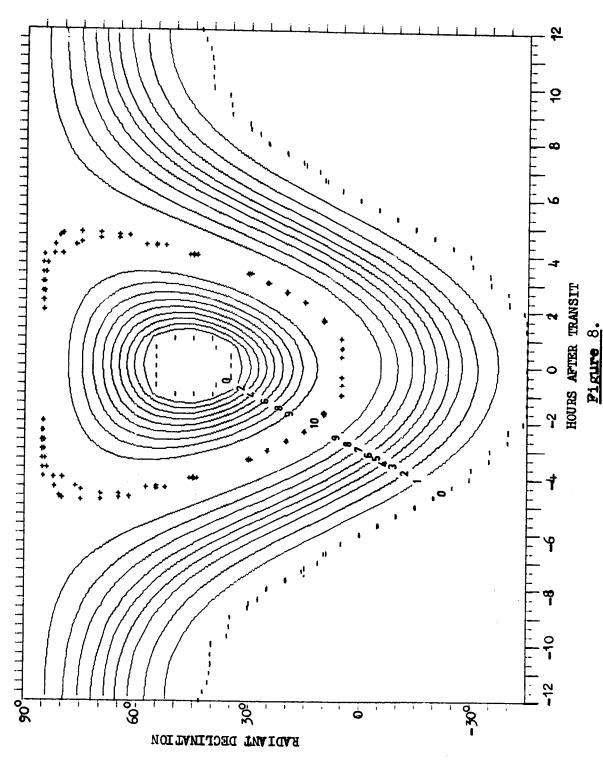
Echo rate as a function of radiant declination and time for the Ottawa system, C = -0.6



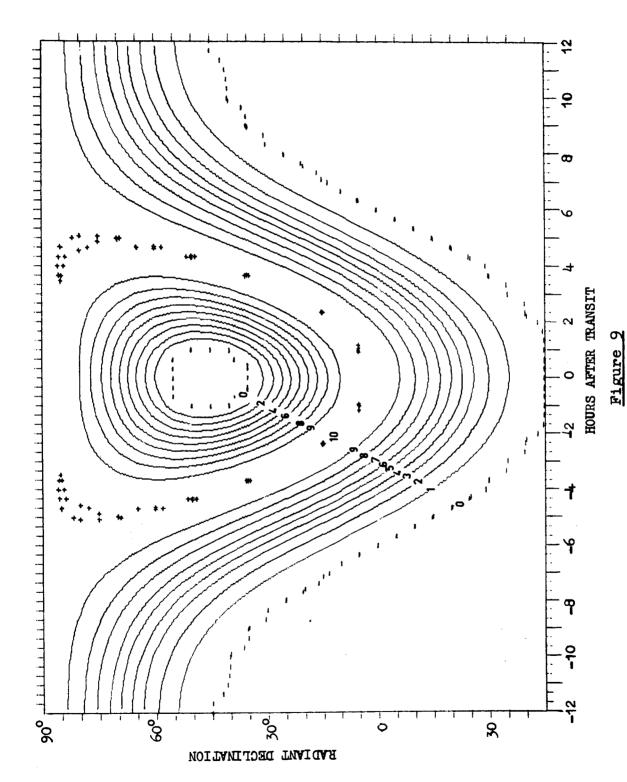
٠. 8 Echo rate as a function of radiant declination and time for Ottawa system, C



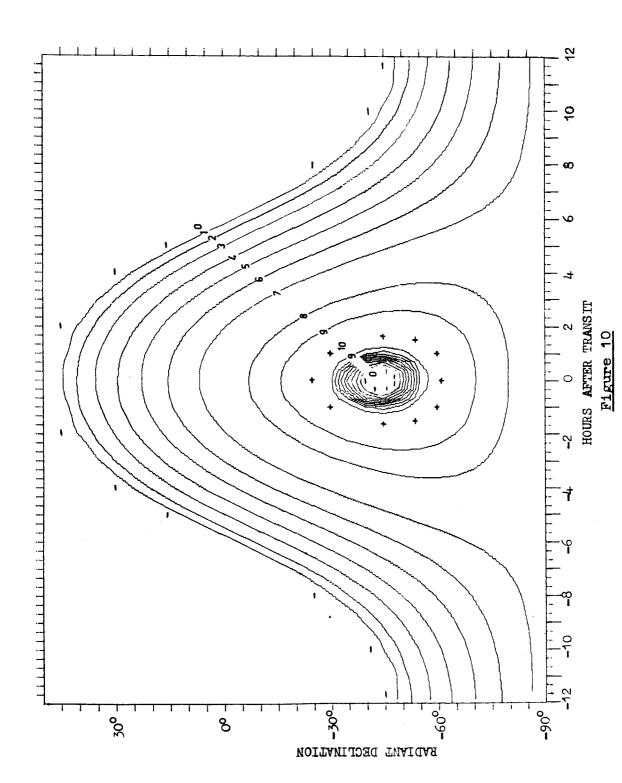
= -1.0 Echo rate as a function of radiant declination and time for the Ottawa system, C



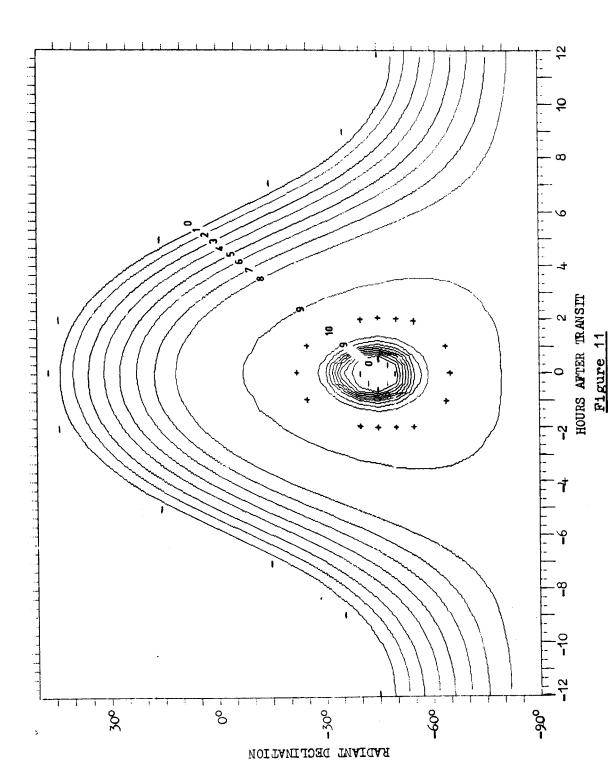
Echo rate as a function of radiant declination and time for the Ottawa system, C = -1.2



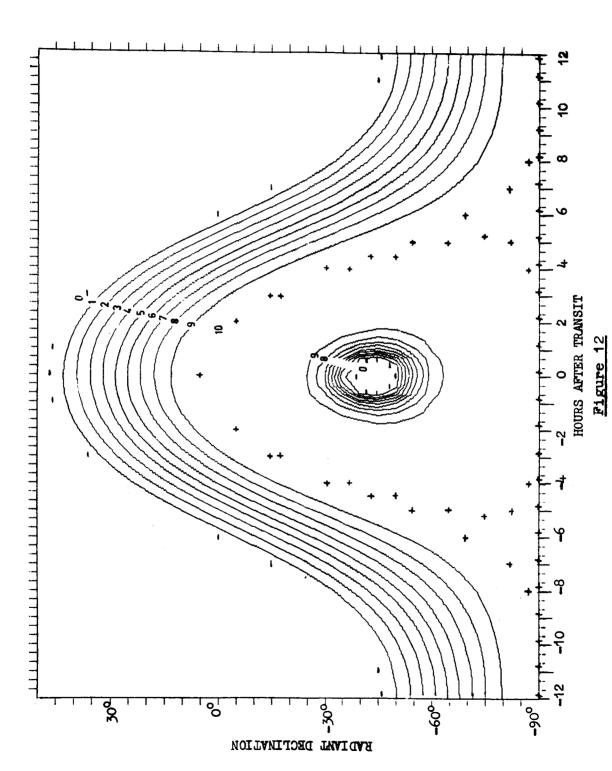
Echo rate as a function of radiant declination and time for the Ottawa system, C = -1.4



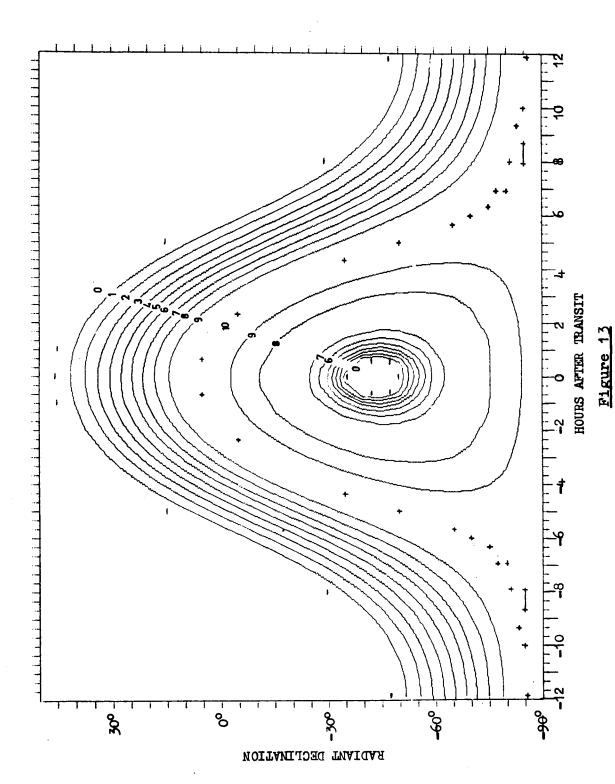
= -0.6 Echo rate as a function of radiant declination and time for the Christchurch system, C



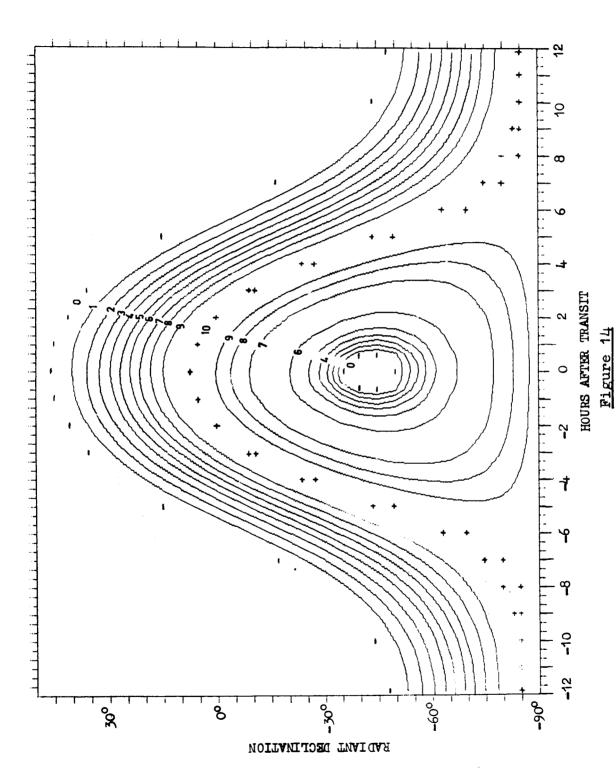
Echo rate as a function of radiant declination and time for the Christchurch system, C=-.8



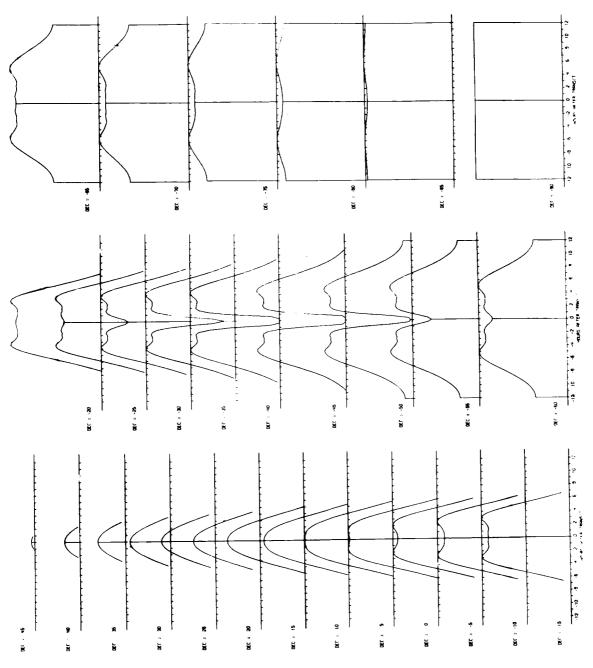
Echo rate as a function of radiant declination and time for the Christdarch system, C =-1.0



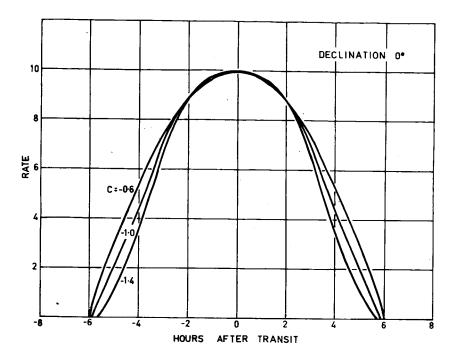
Echo rate as a function of radiant declination and time for the Christchurch system, C = -1.2



Echo rate as a function of radiant declination and time for the Christchurch system, C = -1.4



Diurnal rate curves as a function of radiant declination (5 intervals, C = -1.0). Christchurch. Figure 15.



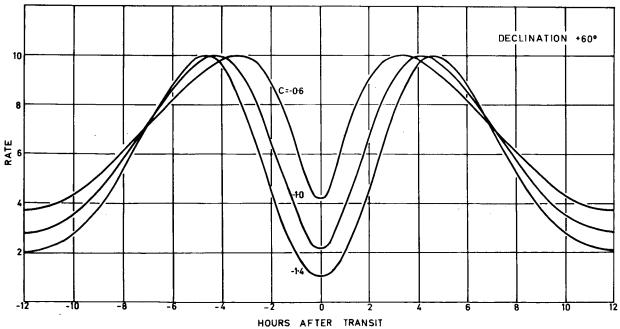
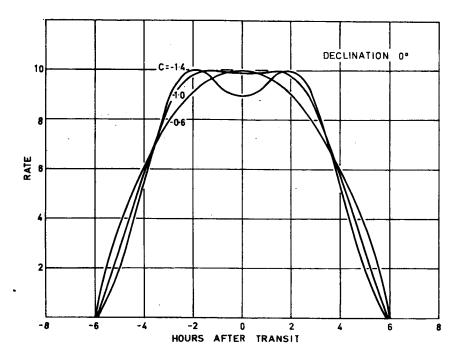


Figure 16. Ottawa diurnal rate curves for radiants at declinations 0° , + 60° and for C = -0.6, -1.0, -1.4. Curves have been normalized to equal maxima.



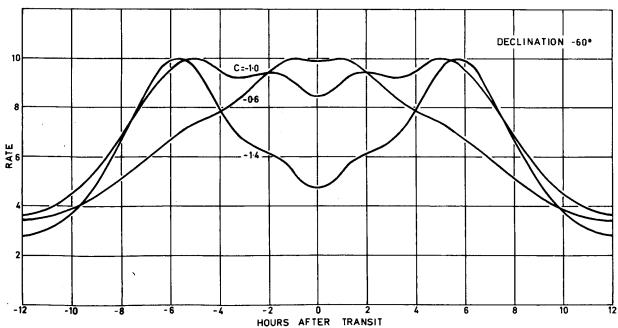


Figure 17. Christchurch diurnal rate curves for radiants at declinations 0° , -60° and for C = -0.6, -1.0, -1.4. Curves have been normalized to equal maxima.